

Structure and Ordering in Metal Cyanide Lattices: the Use of Doubly Labelled Cyanide (^{13}C – ^{15}N) to Simplify the ^{13}C MAS NMR Spectrum†

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The relatively simple procedure of incorporating doubly labelled ^{13}C – ^{15}N into metal cyanide lattices gives a significant simplification of the MAS NMR spectrum, yielding readily interpretable structural information, and the first example of ^{13}C – $^{63,65}\text{Cu}$ J coupling constants in the solid state, for $\text{NMe}_4\text{Cu}(\text{I})\text{Zn}(\text{CN})_4$.

The difficulty in distinguishing C from N by X-ray diffraction has proved to be a serious problem in the study of metal cyanide lattices.^{1–3} Recent ^{113}Cd NMR studies^{4,5} of several cadmium cyanides have shown that ordered lattice space groups chosen on the basis of the refinement of X-ray data were incorrect.

Metal cyanides form an extremely extensive and versatile group of compounds, many with potential as lattice framework hosts. There are few which have a suitable spin 1/2 metal nucleus at the lattice vertices which may serve as NMR probes to check for local order, such as ^{113}Cd in the cadmium cyanides. Hoskins *et al.*¹ recently reported a mixed copper–zinc cyanide, $\text{NMe}_4\text{Cu}(\text{I})\text{Zn}(\text{CN})_4$, and proposed Cu–C–N–Zn ordering, although neither the C and N nor the Cu and Zn atoms could be distinguished explicitly. An initial attempt to use ^{13}C solid-state NMR spectroscopy as an independent method for examining the cyanide ordering resulted in a relatively weak cyanide ^{13}C signal [Fig. 1(a)] which was not easily interpretable. Recognizing that the spectral complexity has its origin in the coupling of ^{13}C to the quadrupolar ^{14}N nucleus, and perhaps also to $^{63,65}\text{Cu}$, we decided to incorpor-

ate doubly labelled ^{13}C – ^{15}N cyanide into $\text{NMe}_4\text{Cu}(\text{I})\text{Zn}(\text{CN})_4$ at the 25% level. The resulting magic angle spinning ^{13}C NMR spectrum [Fig. 1(b, c)] of the enriched sample shows considerable fine structure which yields significant new information by inspection. The spinning sideband manifold consists of quartets which can be assigned to ^{13}C – $^{63,65}\text{Cu}$ J coupling, the first example of its kind. Close inspection of an individual quartet shows partially resolved fine structure on the outer lines, indicative of the slightly different coupling constants between ^{13}C and ^{63}Cu or ^{65}Cu ; $^1J(^{13}\text{C}, ^{63}\text{Cu}) = 324$ and $^1J(^{13}\text{C}, ^{65}\text{Cu}) = 347$ Hz. This observation immediately confirms the Cu–C–N–Zn ordering scheme proposed in the structural study.¹

The ^{13}C NMR spectrum obtained for the same enriched material under static conditions also shows interesting fine structure which results from dipolar couplings to $^{63,65}\text{Cu}$ and ^{15}N , in addition to the J coupling and the ^{13}C chemical shift anisotropy, thus confirming the Cu–C–N–Zn ordering scheme. The ^{13}C static NMR spectrum can be calculated assuming Cu–C–N to be an A–M–X spin system and summing powder pattern contributions from the ($m\text{Cu}, n\text{N}$), $m = -3/2$ to $+3/2$, $n = -1/2$ to $+1/2$, spin states.^{6,7} The calculations are simplified considerably by the colinearity of the two dipolar and the chemical shifts tensors. The ^{15}N NMR spectra obtained under static and spinning conditions offer an

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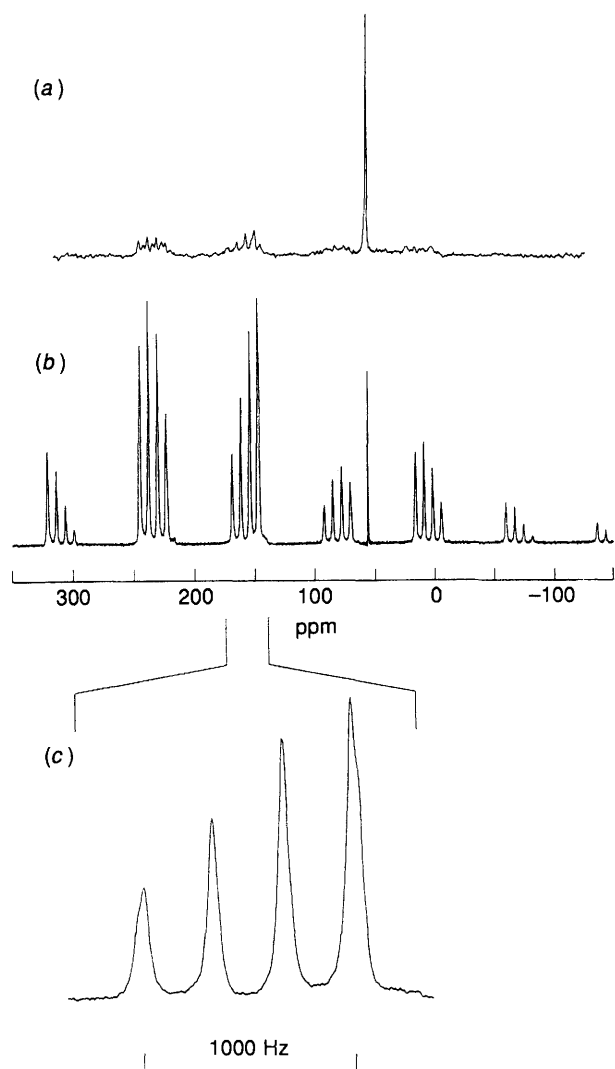


Fig. 1 ^{13}C CPMAS NMR spectra of $\text{NMe}_4\text{Cu}(\text{I})\text{Zn}(\text{CN})_4$: (a) sample with natural isotopic distribution; (b) sample enriched to 25% ^{13}C - ^{15}N cyanide; (c) detail of centrebanded quartet. The line at about 55 ppm is due to the NMe_4^+ carbons.

Table 1 NMR results for doubly labelled $\text{NMe}_4\text{Cu}(\text{I})\text{Zn}(\text{CN})_4$

| Nucleus | Chemical shift anisotropy (ppm) | | | $R(\text{C},\text{N})/\text{Hz}$ | $J(\text{C},\text{N})/\text{Hz}$ | $R(\text{Cu},\text{C})/\text{Hz}$ | $J(\text{Cu},\text{C})/\text{Hz}$ |
|-----------------|---------------------------------|----------------------|-----------------------|----------------------------------|----------------------------------|-----------------------------------|-----------------------------------|
| | δ_{\perp} | δ_{\parallel} | δ_{iso} | | | | |
| ^{15}N | 324 | -105 | 181 | -1788 | -18 | | |
| ^{13}C | 275 | -78 | 157 | -1788 | -18 | 1004 ^a 1075 | 324 ^a 347 |

^a ^{63}Cu values top, ^{65}Cu values bottom.

independent means of evaluating the ^{13}C - ^{15}N dipolar coupling and the ^{13}C - ^{15}N J coupling, 18 Hz, as well as the ^{15}N chemical shift tensor. A calculated ^{13}C powder pattern is shown in Fig. 2, and inspection shows it to be an excellent match for the experimental ^{13}C NMR spectrum. The principal components of the ^{13}C and ^{15}N chemical shift tensors, the dipolar and J couplings, are listed in Table 1.

The dipolar coupling constants R can be converted into internuclear distances r_{am} according to eqn. (1), yielding

$$R = (\mu_0/4\pi) \hbar / 2\pi \gamma_a \gamma_m < r_{\text{am}}^{-3} > \quad (1)$$

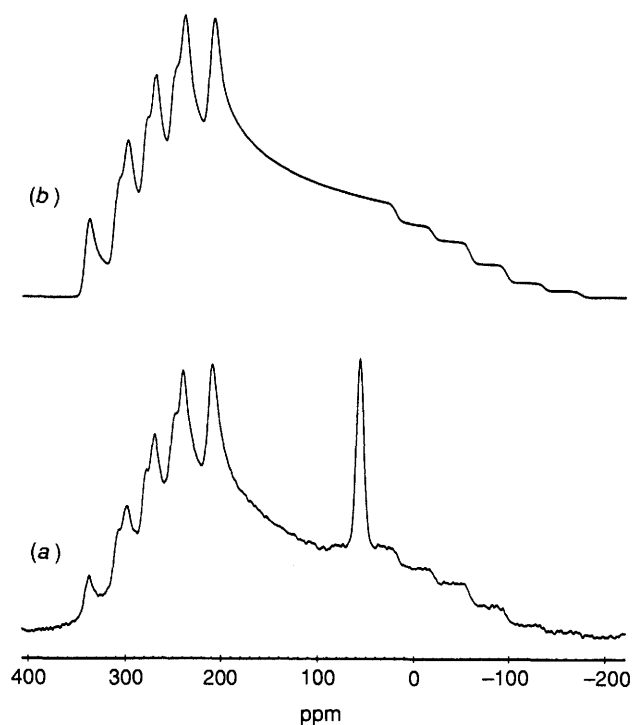


Fig. 2 (a) ^{13}C NMR powder spectrum of $\text{NMe}_4\text{Cu}(\text{I})\text{Zn}(\text{CN})_4$ enriched in doubly labelled cyanide; (b) spectrum calculated with parameters listed in Table 1. The line at about 55 ppm is due to the NMe_4^+ carbons.

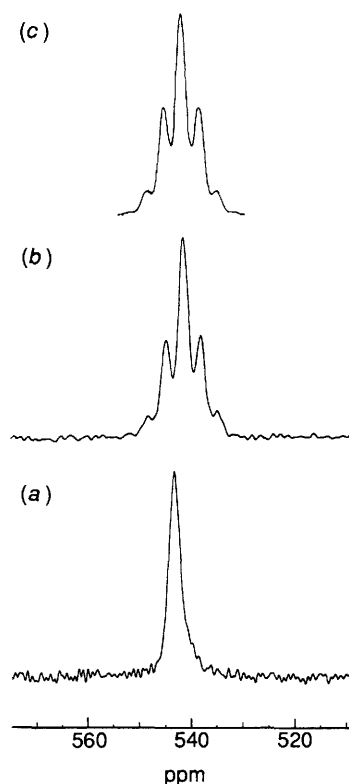


Fig. 3 ^{65}Cu MAS spectra of (a) $\text{NMe}_4\text{Cu}(\text{I})\text{Zn}(\text{CN})_4$; (b) as (a) but sample enriched in doubly labelled cyanide; (c) simulated spectrum assuming a statistical distribution of J couplings to 0, 1, 2 and 3 ^{13}C nuclei appropriate to a ^{13}C incorporation level of 25%.

values of 1.999 and 1.197 Å for Cu-C and C-N. These are considerably longer than the X-ray diffraction values of 1.877 and 1.081 Å respectively. It is important to remember that X-ray diffraction and NMR do not measure exactly the same

quantity, and that both X-ray and NMR values will be modified considerably by the presence of motional averaging. In fact the unusual thermal ellipsoids of the C and N atoms obtained from the analysis of the X-ray diffraction results are indicative of some motion of the CN groups. Also, there is an unknown contribution from anisotropic J coupling, so that the effective dipolar coupling is $R_{\text{eff}} = S(R - \Delta J/3)$ with $\Delta J = J_{\parallel} - J_{\perp}$ and S is an order parameter which depends on the motion. A more detailed analysis will depend on ^{13}C and ^{15}N static NMR spectra obtained at low temperatures so as to reduce the effect of motional averaging.

Finally, we present a ^{65}Cu NMR spectrum for both the unlabelled and doubly labelled samples of $\text{NMe}_4\text{-Cu(I)Zn(CN)}_4$ (Fig. 3). The NMR spectrum of the labelled material shows superimposed resonances of ^{65}Cu attached to 0, 1, 2 and 3 ^{13}C nuclei, with corresponding multiplets arising from the J coupling. The relative intensities of the multiplets follow a statistical distribution appropriate for 25% labelling. The relatively sharp ^{65}Cu spectrum, owing to the presence of Cu on exact tetrahedral lattice sites, makes the natural abundance material suitable as a chemical shift and cross-polarization standard for solid-state Cu NMR studies.

Incorporation of ^{13}C , ^{15}N doubly enriched cyanide has a tremendous potential for elucidating local order and structure

in metal cyanides. The incorporation is usually relatively simple, and at enrichment levels of 10–25% the natural abundance ^{14}N -coupled ^{13}C NMR signals become essentially invisible.

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